

A 12 PART WEBINAR SERIES ON:

# COMPOSITE MATERIALS ENGINEERING

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Co-Director, Master of Engineering Leadership, AMM Program, UBC  
Lead of Continuing Professional Development, CKN

- Ph.D. and M.A.Sc. in Composite Materials Engineering
- Over 15 years experience in industry and academia working on polymer matrix composites in aerospace, automotive, marine, energy, recreation and others
- Experience working with over 150 companies from SME to major international corporations
- Expertise in liquid composite moulding and thermal management

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WEBINARS

## REIGNITING THE CANADIAN COMPOSITES MANUFACTURING INDUSTRY

12-PART WEBINAR SERIES



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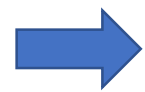
# OVERVIEW OF WEBINAR SERIES

- Series of 12 webinars, 1 hour each

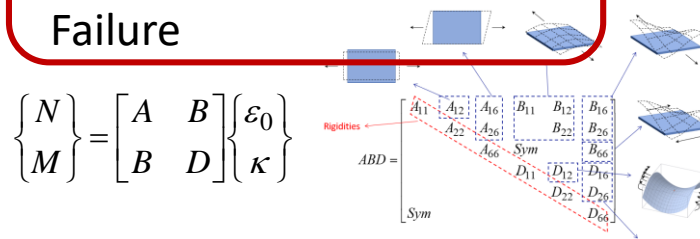
Introduction  
Constituent Materials  
Thermal Management



Processing (Manufacturing)  
Prepreg Processing  
Liquid Composite Moulding



Mechanics of Composites  
Micromechanics  
Lamina and Laminate Level  
Failure



Testing Composites  
Common Defects



- For more information on dates and times visit:

<https://compositeskn.org/aimevents/>

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## OUTCOMES OF THIS WEBINAR

- How can I predict my material properties based on my selection of resin and fibre?
- What effect does  $V_f$  have on material properties?
- What effect does fibre orientation have on material properties?
- **Upcoming webinars**
  - What effect does fibre, matrix, and stacking sequence have on material properties?
  - What effect does fibre, matrix, and stacking sequence have on strength?

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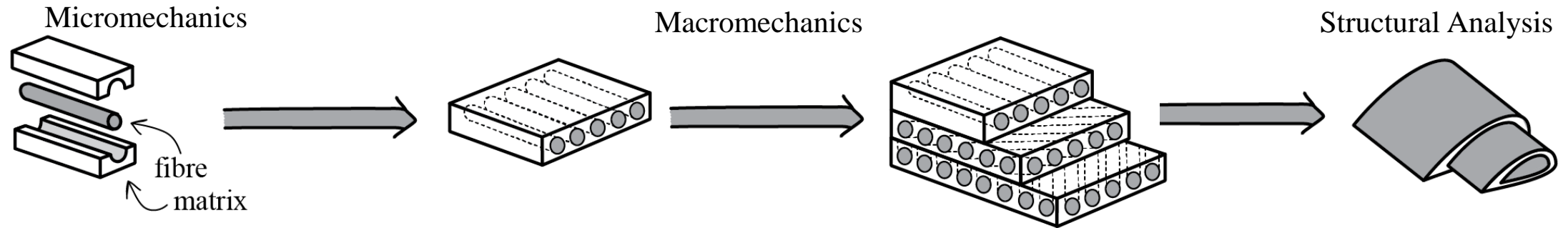


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# ANALYSIS OF COMPOSITE MATERIALS

The complicated nature of a structure made from composite materials lends to many levels of analysis



- Micro-mechanics focuses on the stresses, strains, damage, etc., happening at the level of individual fibres and the fibre/matrix interface
- Macro-mechanics focuses on the analysis at the level of individual lamina and sub-laminates
- Structural mechanics involves analysis at the structural scale

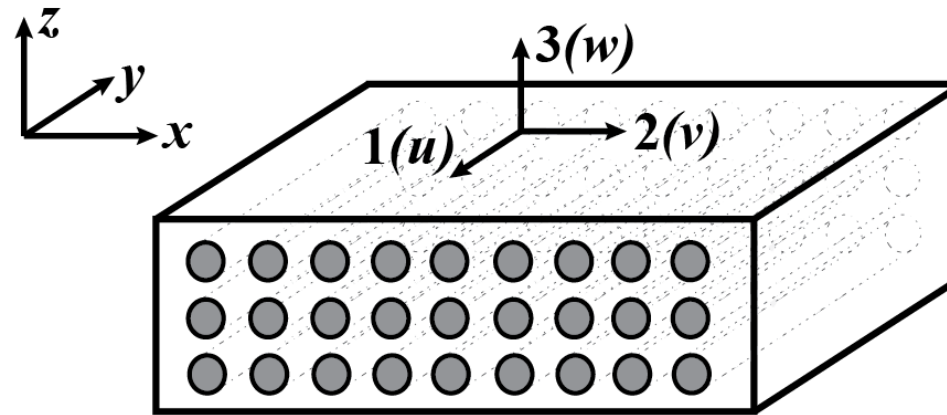
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## CONVENTIONS

- A common 'dialect' has been developed among those who perform design and analysis with composite materials
- The fibre direction in a lamina (or a single layer) is called the "1" direction
  - "u" = displacement in 1-direction
- The in-plane perpendicular to the fibre direction is called the "2" direction
  - "v" = displacement in 2-direction
- The out-of-plane perpendicular to the fibre direction is called the "3" direction
  - "w" = displacement in 3-direction

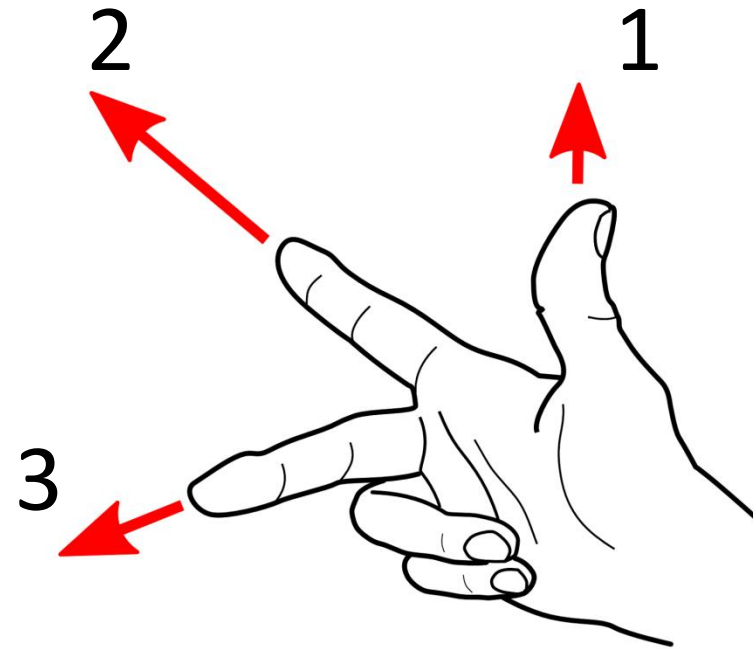
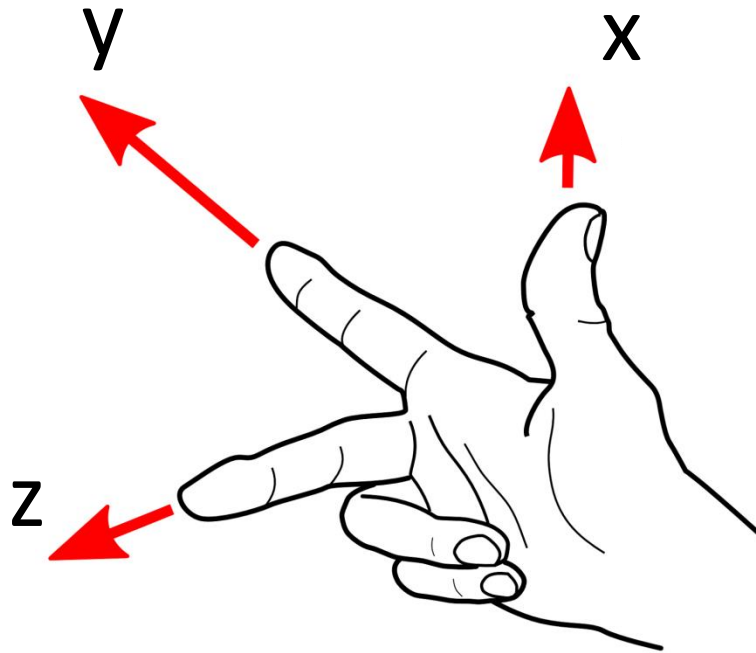


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# CONVENTIONS

- Right-hand rule



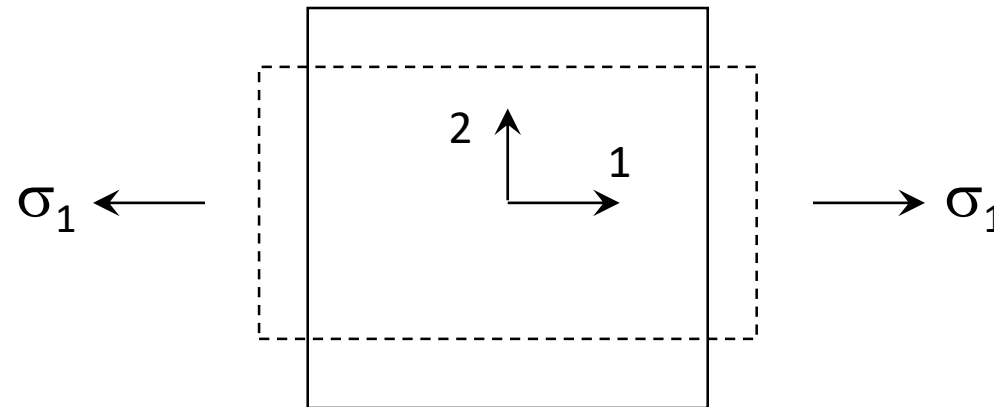
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## HOOKE'S LAW: 1-DIRECTION

- Consider a linear elastic isotropic material (e.g. steel). It is subjected only to stresses in the 1- or 2- directions. This is known as *plane stress*. Plane stress is an assumption that is valid when the thickness of a body is small and there are no out-of-plane loads, typical conditions for a lamina.
- Now consider such a material subjected to an axial load:



## HOOKE'S LAW: 1-DIRECTION

- The resulting strains can be written as:

$$\varepsilon_1 = \frac{\sigma_1}{E} \qquad \varepsilon_2 = -\nu\varepsilon_1 = -\nu \frac{\sigma_1}{E}$$

- Where  $E$  is the Young's modulus (modulus of elasticity) and  $\nu$  is the Poisson's ratio. Recall that Poisson's ratio is defined as:

$$\nu = -\frac{\varepsilon_2}{\varepsilon_1}$$

under uniaxial loading conditions.

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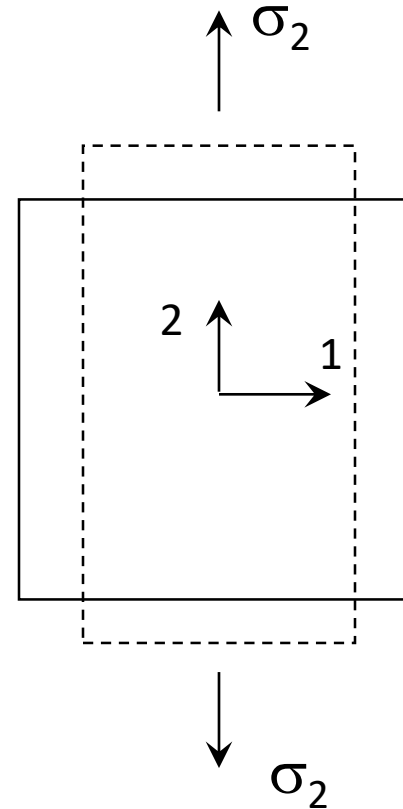
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## HOOKE'S LAW: 2-DIRECTION

- Similarly, if a uniaxial load is considered in the 2-direction:

$$\varepsilon_1 = -\nu \varepsilon_2 = -\nu \frac{\sigma_2}{E}$$

$$\varepsilon_2 = \frac{\sigma_2}{E}$$



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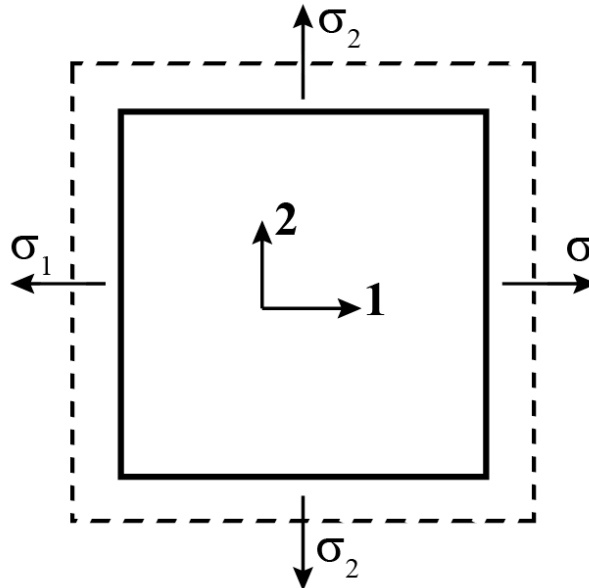
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## HOOKE'S LAW: BIAXIAL LOAD

- If both stresses are present, then we can express the strains by using the principle of superposition

$$\varepsilon_1 = \frac{\sigma_1}{E} - \nu \frac{\sigma_2}{E}$$

$$\varepsilon_2 = -\nu \frac{\sigma_1}{E} + \frac{\sigma_2}{E}$$



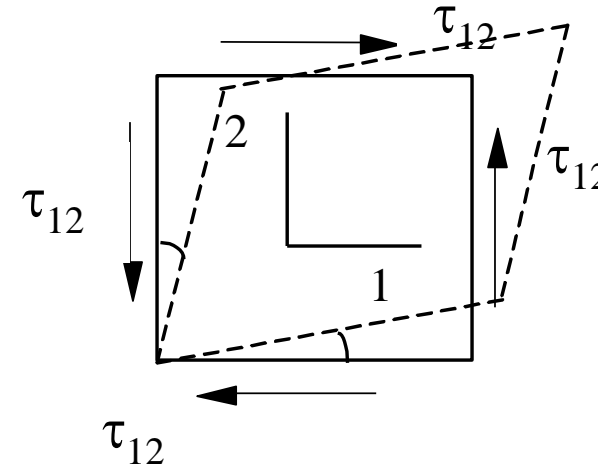
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## HOOKE'S LAW: SHEAR STRESS

- In a similar manner, consider the case where only a shear stress is applied

$$\gamma_{12} = \frac{\tau_{12}}{G}$$



- $G$ , the shear modulus, can be expressed in terms of  $E$  and  $\nu$  for isotropic materials (ie. only two independent elastic constants):

$$G = \frac{E}{2(1 + \nu)}$$

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## HOOKE'S LAW

- The biaxial and pure shear loading cases can be combined and written in matrix form for an isotropic material:

$$\begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix} = \underbrace{\begin{bmatrix} \frac{1}{E} & \frac{-\nu}{E} & 0 \\ \frac{-\nu}{E} & \frac{1}{E} & 0 \\ 0 & 0 & \frac{1}{G} \end{bmatrix}}_{\text{Compliance Matrix [S]}} \begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix} \quad \begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix} = \underbrace{\begin{bmatrix} \frac{E}{1-\nu^2} & \frac{\nu E}{1-\nu^2} & 0 \\ \frac{\nu E}{1-\nu^2} & \frac{E}{1-\nu^2} & 0 \\ 0 & 0 & G \end{bmatrix}}_{\text{Stiffness Matrix [Q]}} \begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix}$$

Compliance Matrix [S]

Stiffness Matrix [Q]

$$\{\varepsilon\} = [S]\{\sigma\}$$

$$[Q] = [S]^{-1}$$

$$\{\sigma\} = [Q]\{\varepsilon\}$$

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## ORTHOTROPIC MATERIALS

- Composite materials are orthotropic, however, the same principles as isotropic apply.
- In the 1-direction, the strain equations are:

$$\varepsilon_1 = \frac{\sigma_1}{E_1}$$

$$\varepsilon_2 = -\nu_{12}\varepsilon_1 = -\nu_{12} \frac{\sigma_1}{E_1}$$

- and, in the 2-direction:

$$\varepsilon_2 = \frac{\sigma_2}{E_2}$$

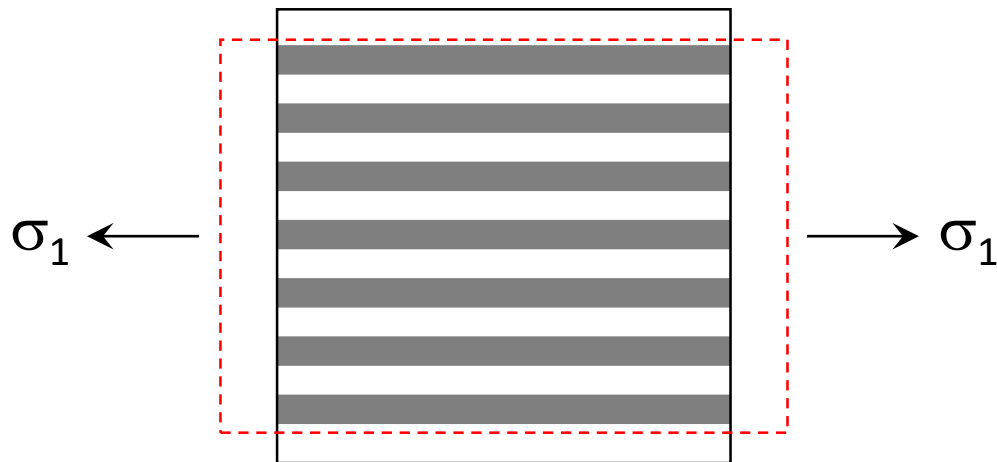
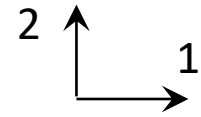
$$\varepsilon_1 = -\nu_{21}\varepsilon_2 = -\nu_{21} \frac{\sigma_2}{E_2}$$

- Note that for unidirectional composites,  $E_1 \gg E_2$ , and typically  $\nu_{12} > \nu_{21}$ . As a result,  $\nu_{12}$  is the “major” Poisson’s ratio,  $\nu_{21}$  is the “minor” Poisson’s ratio.

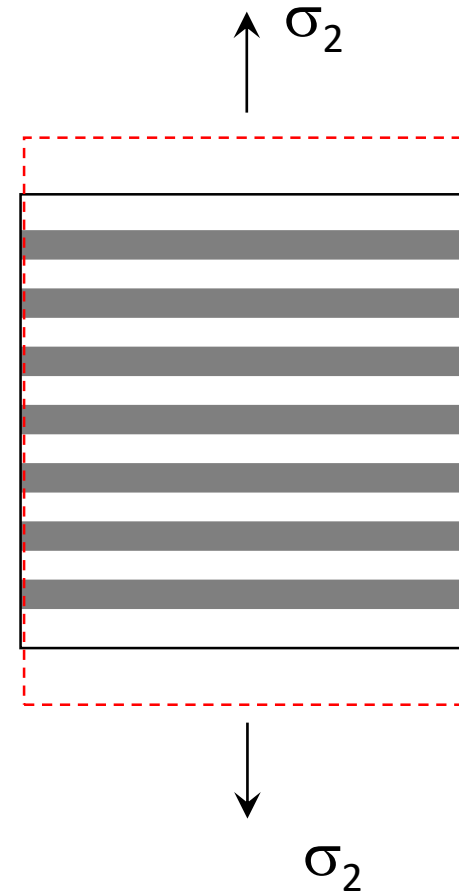
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# ORTHOTROPIC MATERIALS: POISSON'S RATIO



Effect of Major Poisson's Ratio,  $\nu_{12}$



Effect of Minor Poisson's Ratio,  $\nu_{21}$

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## HOOKE'S LAW FOR ORTHOTROPIC MATERIALS

- The stress-strain relationship can be written:

$$\begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix} = \begin{bmatrix} \frac{1}{E_1} & \frac{-\nu_{21}}{E_2} & 0 \\ \frac{-\nu_{12}}{E_1} & \frac{1}{E_2} & 0 \\ 0 & 0 & \frac{1}{G_{12}} \end{bmatrix} \begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix}$$

- However, for any elastic body, the compliance and stiffness matrices must be symmetric. In this case:

$$\frac{\nu_{21}}{E_2} = \frac{\nu_{12}}{E_1}$$

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# HOOKE'S LAW FOR ORTHOTROPIC MATERIALS

- Hooke's Law for a 2D unidirectional lamina is then:

$$\begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix} = \underbrace{\begin{bmatrix} \frac{1}{E_1} & -\frac{\nu_{12}}{E_1} & 0 \\ -\frac{\nu_{12}}{E_1} & \frac{1}{E_2} & 0 \\ 0 & 0 & \frac{1}{G_{12}} \end{bmatrix}}_{\text{Compliance Matrix [S]}} \begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix}$$

Compliance Matrix [S]

$$\begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix} = \underbrace{\begin{bmatrix} \frac{E_1}{1-\nu_{21}\nu_{12}} & \frac{\nu_{12}E_2}{1-\nu_{21}\nu_{12}} & 0 \\ \frac{\nu_{12}E_2}{1-\nu_{21}\nu_{12}} & \frac{E_2}{1-\nu_{21}\nu_{12}} & 0 \\ 0 & 0 & G_{12} \end{bmatrix}}_{\text{Stiffness Matrix [Q]}} \begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix}$$

Stiffness Matrix [Q]

- Note that 4 independent constants are required:

$$E_1, E_2, (\nu_{12} \text{ or } \nu_{21}), G_{12}$$

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## MICROMECHANICS

- Given a lamina, made from a specific matrix and a specific fibre, how are the engineering constants determined?
- The engineering constants depend on a number of variables:
  - The individual material constituents
  - The fibre volume fraction
  - The packing geometry
  - The processing method

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## MICROMECHANICS: VOLUME FRACTION

- $V_{c,f,m,v}$  = volume of composite, fiber, matrix, voids
- $\rho_{c,f,m}$  = density of composite, fiber, matrix
- The fiber volume fraction,  $V_f$  is:

$$V_f = \frac{V_f}{V_c}$$

- The matrix volume fraction,  $V_m$  is:

$$V_m = \frac{V_m}{V_c}$$

- The void volume fraction,  $V_v$  is:

$$V_v = \frac{V_v}{V_c}$$

- Note that:

$$V_f + V_m + V_v = 1$$

$$V_f + V_m + V_v = V_c$$

*Note:  $V_*$  = volume fraction (big V)  
 $v_*$  = volume of component (little v)*

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## MICROMECHANICS: MASS FRACTION

- $w_{c,f,m}$  = mass of composite, fiber, matrix
- The fiber mass fraction,  $W_f$  is:

$$W_f = \frac{W_f}{W_c}$$

- The matrix mass fraction,  $w_m$  is:

$$W_m = \frac{W_m}{W_c}$$

- The total mass of the composite,  $w_c$  is:

$$W_c = W_f + W_m$$

*Note:  $W_*$  = mass fraction (big W)  
 $w_*$  = mass of component (little w)*

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## MICROMECHANICS: DENSITY

- Knowing the volume and masses, we can determine the density:

$$W_c = W_f + W_m$$

$$\rho_c v_c = \rho_f v_f + \rho_m v_m$$

$$\rho_c = \rho_f \frac{v_f}{v_c} + \rho_m \frac{v_m}{v_c}$$

$$\rho_c = \rho_f V_f + \rho_m V_m$$

Assuming voids are negligible:

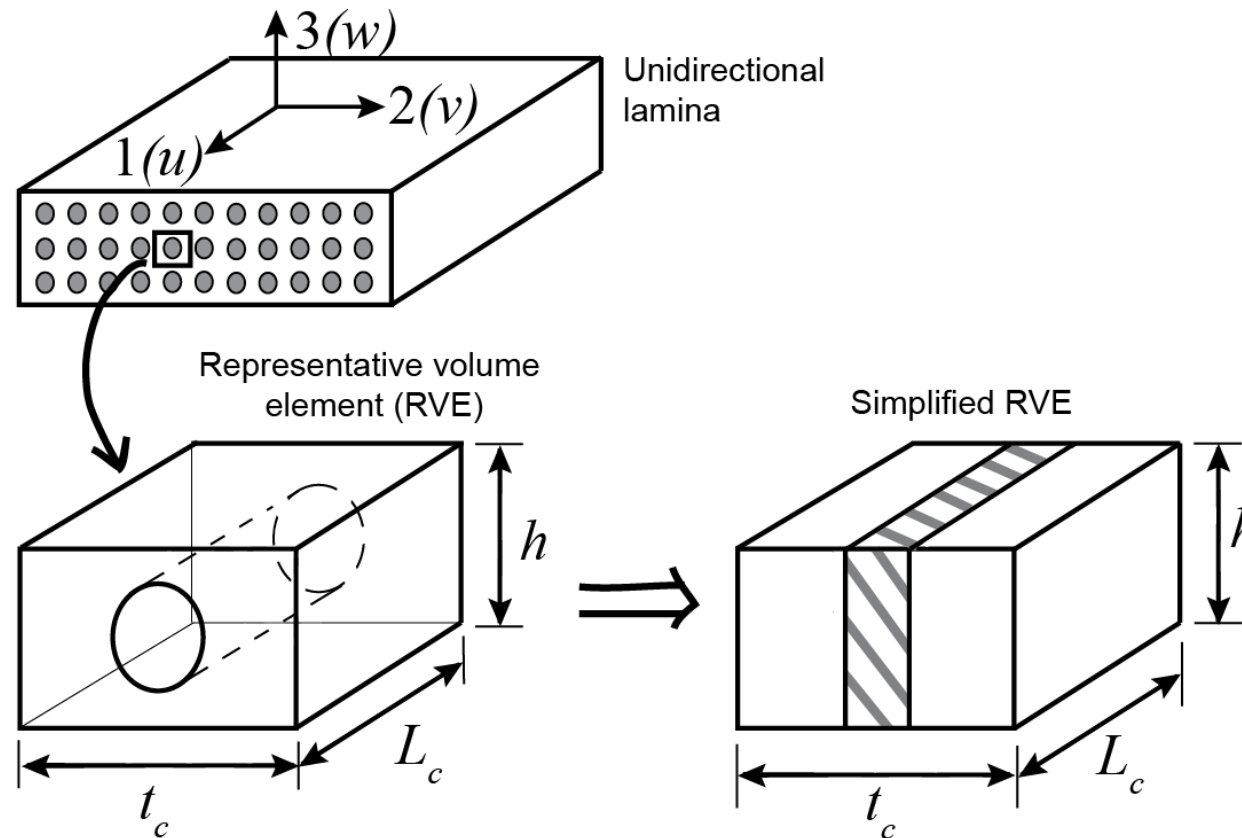
$$\rho_c = \rho_f V_f + \rho_m (1 - V_f)$$

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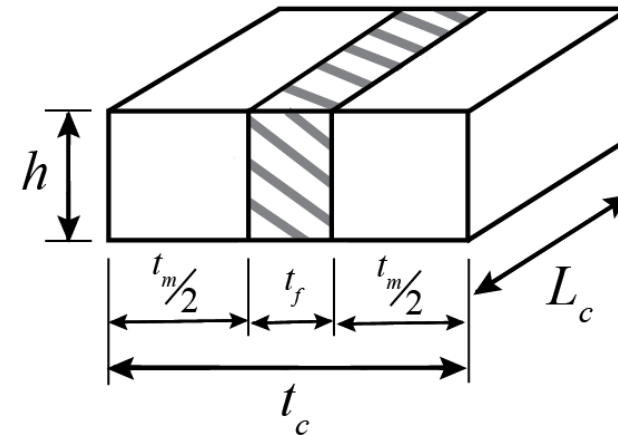
# MICROMECHANICS: STRENGTH OF MATERIALS APPROACH

- Consider a representative volume element (RVE) taken from a unidirectional lamina:



# MICROMECHANICS: SIMPLIFIED RVE REPRESENTATION

- The fibre, matrix and composite are assumed to be the same height,  $h$  but of thicknesses  $t_f$ ,  $t_m$ , and  $t_c$
- The area of each are:



$$A_f = t_f h$$

$$A_m = t_m h$$

$$A_c = t_c h$$

- The volume fractions are defined as:

$$V_f = \frac{A_f}{A_c} = \frac{t_f}{t_c}$$

$$V_m = \frac{A_m}{A_c} = \frac{t_m}{t_c} = 1 - V_f$$

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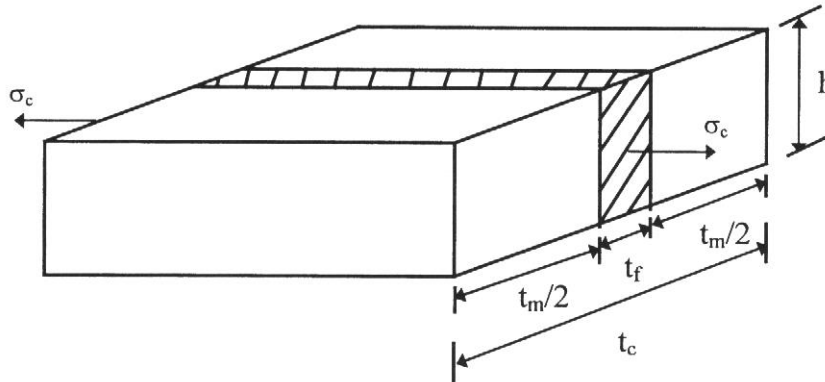
## MICROMECHANICS: STRENGTH OF MATERIALS APPROACH

- The strength of materials approach (also known as the mechanics of materials approach) to micro-mechanics makes the following assumptions:
  - The bond between the fibres and matrix is perfect
  - The fibre diameters, and the space between the fibres are uniform
  - The fibres are continuous and parallel
  - The fibres and matrix are linear elastic, following Hooke's Law
  - The fibres are all the same strength and the elastic moduli of the fibres and matrix are constant
  - The composite contains no voids

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# MICROMECHANICS: LONGITUDINAL YOUNG'S MODULUS



- Consider a unidirectional force applied to the simplified RVE in the fibre direction, such that it is shared by the fibre and the matrix

$$F_c = F_f + F_m$$

- These forces can also be expressed as stresses:

$$F_c = \sigma_c A_c$$

$$F_f = \sigma_f A_f$$

$$F_m = \sigma_m A_m$$

- With the assumption of linear-elastic, isotropic fibres and matrix we can write the stress-strain relationships as:

$$\sigma_c = E_1 \varepsilon_c$$

$$\sigma_f = E_f \varepsilon_f$$

$$\sigma_m = E_m \varepsilon_m$$

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# MICROMECHANICS: LONGITUDINAL YOUNG'S MODULUS

- Now, we can write:

$$F_c = F_f + F_m$$

$$E_1 \varepsilon_c A_c = E_f \varepsilon_f A_f + E_m \varepsilon_m A_m$$

- But, our assumptions tell us:

$$\delta_c = \delta_f = \delta_m$$

- And since the lengths are the same:

$$\varepsilon_c = \varepsilon_f = \varepsilon_m$$

- So, we can write:

$$E_1 = E_f \frac{A_f}{A_c} + E_m \frac{A_m}{A_c}$$

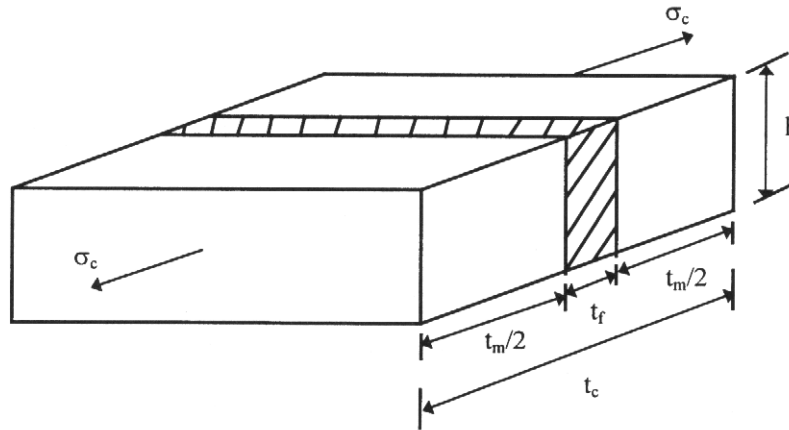
- Or, recalling the volume fractions:

$$E_1 = E_f V_f + E_m V_m$$

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# MICROMECHANICS: TRANSVERSE YOUNG'S MODULUS



- Now consider a unidirectional force applied to the simplified RVE in the matrix direction

- In this case, the force in the matrix and the fibre is the same, but the displacements add up to the total displacement:

$$F_c = F_f = F_m$$

$$\sigma_c = \sigma_f = \sigma_m$$

$$\delta_c = \delta_f + \delta_m$$

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## MICROMECHANICS: TRANSVERSE YOUNG'S MODULUS

- Unlike the longitudinal case, the thickness of the matrix and fibre are different, so the strains are not equal:

$$\delta_c = t_c \varepsilon_c$$

$$\delta_f = t_f \varepsilon_f$$

$$\delta_m = t_m \varepsilon_m$$

- But, we can express the strains according to Hooke's Law:

$$\varepsilon_c = \frac{\sigma_c}{E_2}$$

$$\varepsilon_f = \frac{\sigma_f}{E_f}$$

$$\varepsilon_m = \frac{\sigma_m}{E_m}$$

- Therefore, the displacement equation is:

$$t_c \frac{\sigma_c}{E_2} = t_f \frac{\sigma_f}{E_f} + t_m \frac{\sigma_m}{E_m}$$

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## MICROMECHANICS: TRANSVERSE YOUNG'S MODULUS

- However, the stresses are the same and:

$$\frac{t_f}{t_c} = \frac{t_f h}{t_c h} = \frac{A_f}{A_c} = V_f \qquad \frac{t_m}{t_c} = \frac{t_m h}{t_c h} = \frac{A_m}{A_c} = V_m$$

- Therefore:

$$t_c \frac{\sigma_c}{E_2} = t_f \frac{\sigma_f}{E_f} + t_m \frac{\sigma_m}{E_m}$$

$$\frac{1}{E_2} = \frac{1}{E_f} V_f + \frac{1}{E_m} V_m$$

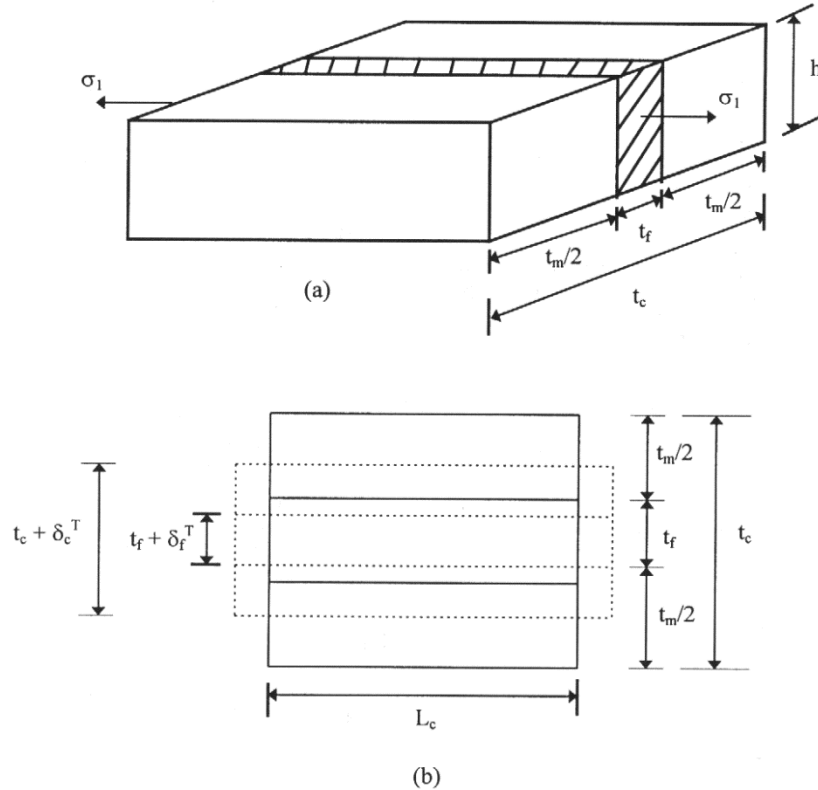
- Rearranged:

$$E_2 = \frac{E_f E_m}{E_m V_f + E_f V_m}$$

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# MICROMECHANICS: MAJOR POISSON'S RATIO



- Assume our RVE is loaded in the 1-direction (along the fibres)
- The *transverse* displacements are:

$$\delta_c^T = \delta_f^T + \delta_m^T$$

- The *transverse* strains are:

$$\varepsilon_c^T = \frac{\delta_c^T}{t_c} \quad \varepsilon_f^T = \frac{\delta_f^T}{t_f} \quad \varepsilon_m^T = \frac{\delta_m^T}{t_m}$$

- And, from the definition of Poisson's ratio:

$$\varepsilon_c^T = -\nu_{12} \varepsilon_c^L \quad \nu_{LT} = \frac{-\varepsilon^T}{\varepsilon^L}$$

$$\varepsilon_f^T = -\nu_f \varepsilon_f^L$$

$$\varepsilon_m^T = -\nu_m \varepsilon_m^L$$

*Note: Poisson's ratio is the Greek letter  $\nu$  (nu), not the letter  $v$*

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## MICROMECHANICS: MAJOR POISSON'S RATIO

- So:

$$\delta_c^T = \delta_f^T + \delta_m^T$$

$$-t_c \nu_{12} \varepsilon_c^L = -t_f \nu_f \varepsilon_f^L - t_m \nu_m \varepsilon_m^L$$

- But, when loading in the 1-direction, the strains are equal...

$$\nu_{12} = \nu_f \frac{t_f}{t_c} + \nu_m \frac{t_m}{t_c}$$

- And, recalling the volume fraction equations:

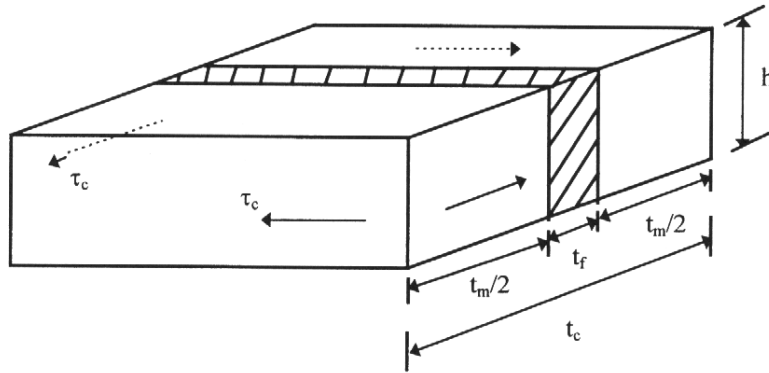
$$\nu_{12} = \nu_f V_f + \nu_m V_m$$

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# MICROMECHANICS: IN-PLANE SHEAR MODULUS



- Consider a pure shear stress,  $\tau_c$ , applied to the RVE
- The resulting deformations are:

$$\delta_c = \delta_f + \delta_m$$

where the displacements are:

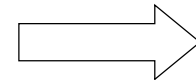
$$\delta_c = \gamma_c t_c \quad \delta_f = \gamma_f t_f \quad \delta_m = \gamma_m t_m$$

- Now, recalling Hooke's Law,

$$\gamma_c = \frac{\tau_c}{G_{12}} \quad \gamma_f = \frac{\tau_f}{G_f} \quad \gamma_m = \frac{\tau_m}{G_m}$$

therefore:

$$\frac{\tau_c}{G_{12}} t_c = \frac{\tau_f}{G_f} t_f + \frac{\tau_m}{G_m} t_m$$



$$\frac{1}{G_{12}} = \frac{1}{G_f} \frac{t_f}{t_c} + \frac{1}{G_m} \frac{t_m}{t_c}$$

$$\frac{1}{G_{12}} = \frac{1}{G_f} V_f + \frac{1}{G_m} V_m$$

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## MICROMECHANICS: SUMMARY

$$\rho_c = \rho_f V_f + \rho_m V_m$$

$$E_1 = E_f V_f + E_m V_m$$

$$\nu_{12} = \nu_f V_f + \nu_m V_m$$

$$\frac{1}{E_2} = \frac{1}{E_f} V_f + \frac{1}{E_m} V_m$$

$$\frac{1}{G_{12}} = \frac{1}{G_f} V_f + \frac{1}{G_m} V_m$$

AKA “Rule of Mixtures”

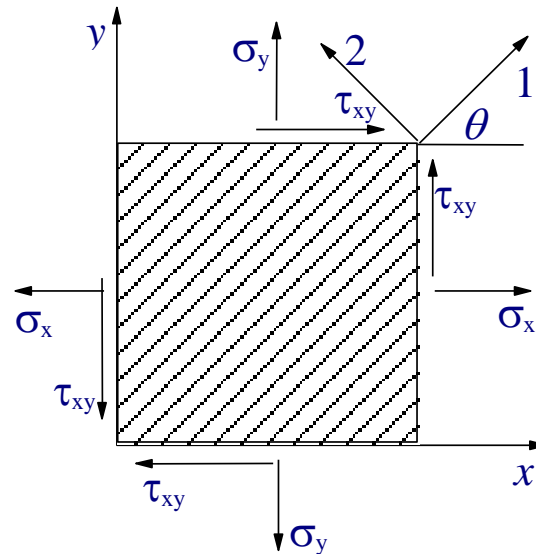
$$X = X_1 V_1 + X_2 V_2 + \cdots + X_n V_n$$

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## ANGLE LAMINA

- When the material axes (1-2-3) of a lamina do not coincide with the loading axes (x-y-z), then the lamina is an “angle lamina”
- In order to determine the stress-strain relations in the global coordinate system (x-y-z), the elastic properties of the lamina have to be transformed from the local coordinate system (1-2-3)



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## (STRESS) TRANSFORMATION MATRIX

- Using a transformation matrix, stresses along the x-y directions can be related to the stresses in the 1-2 directions:

$$\begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix} = \begin{bmatrix} c^2 & s^2 & 2sc \\ s^2 & c^2 & -2sc \\ -sc & sc & c^2 - s^2 \end{bmatrix} \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix}$$

where  $s = \sin\theta$  , and  $c = \cos\theta$

- This can be written as:

$$\begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix} = [T] \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix}$$

## (STRAIN) TRANSFORMATION MATRIX

- Strains are transformed in the same manner:

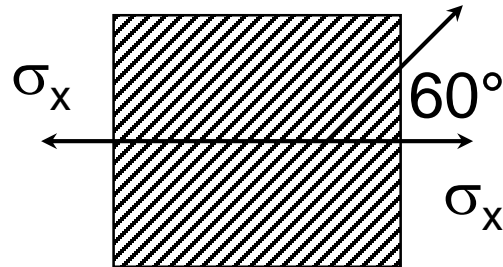
$$\begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix} = [T^*] \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix}$$

- where  $[T^*]$  is

$$[T^*] = \begin{bmatrix} c^2 & s^2 & sc \\ s^2 & c^2 & -sc \\ -2sc & 2sc & c^2 - s^2 \end{bmatrix}$$

## STRESS/STRAIN TRANSFORMATION EXAMPLE

- A unidirectional lamina is oriented at  $60^\circ$ . A stress of 10 MPa is applied in the x-direction ( $0^\circ$ )
  - Determine the stresses in the material direction (1-2)



$$[T] = \begin{bmatrix} c^2 & s^2 & 2sc \\ s^2 & c^2 & -2sc \\ -sc & sc & c^2 - s^2 \end{bmatrix} = \begin{bmatrix} \frac{1}{4} & \frac{3}{4} & \frac{\sqrt{3}}{2} \\ \frac{3}{4} & \frac{1}{4} & -\frac{\sqrt{3}}{2} \\ -\frac{\sqrt{3}}{4} & \frac{\sqrt{3}}{4} & -\frac{1}{2} \end{bmatrix}$$

$$\begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix} = [T] \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = [T] \begin{Bmatrix} 10 \\ 0 \\ 0 \end{Bmatrix} = \begin{Bmatrix} 2.5 \\ 7.5 \\ -4.33 \end{Bmatrix} \text{ MPa}$$

## STRESS/STRAIN TRANSFORMATION

- Conversion to the x-y-z coordinate system works in the same manner:

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = [T]^{-1} \begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix}$$

$$[T]^{-1} = \begin{bmatrix} c^2 & s^2 & -2sc \\ s^2 & c^2 & 2sc \\ sc & -sc & c^2 - s^2 \end{bmatrix}$$

## HOOKE'S LAW FOR ANGLE LAMINA

- Recall Hooke's Law:

$$\begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix} = [S] \begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix}$$

- We need to establish a similar relationship:

$$\begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = [\bar{S}] \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix}$$

Note:  $[\bar{S}]$  is referred to as the transformed compliance matrix

- First, replace the stress and strain terms:

$$[T^*] \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = [S][T] \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix}$$

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## HOOKE'S LAW FOR ANGLE LAMINA

- Now, pre-multiply by the inverse of  $[T^*]$

$$\begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = [T^*]^{-1} [S][T] \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix}$$

- So,

$$[\bar{S}] = [T^*]^{-1} [S][T]$$

- Now,  $[S]$  is a compliance matrix in the x-y-z system:

$$[\bar{S}] = \begin{bmatrix} \bar{S}_{11} & \bar{S}_{12} & \bar{S}_{16} \\ & \bar{S}_{22} & \bar{S}_{26} \\ Sym & & \bar{S}_{66} \end{bmatrix}$$

Note:  $[\bar{S}]$  is referred to as the transformed compliance matrix

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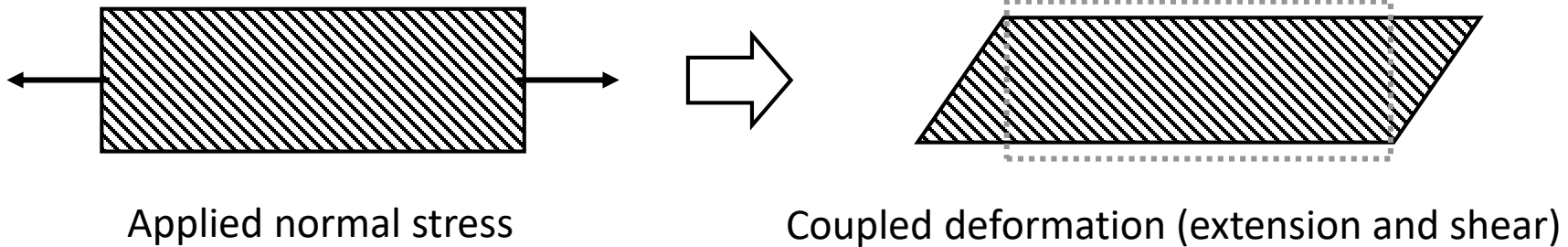
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# COMPLIANCE MATRIX FOR ANGLE LAMINA

$$[\bar{S}] = \begin{bmatrix} \bar{S}_{11} & \bar{S}_{12} & \bar{S}_{16} \\ & \bar{S}_{22} & \bar{S}_{26} \\ Sym & & \bar{S}_{66} \end{bmatrix}$$

Non-zero terms in an angle lamina means that there is coupling between shear and axial directions

What does this mean in reality?



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## ENGINEERING CONSTANTS FOR ANGLE LAMINA

- Based on the compliance matrix of an angle lamina, engineering constants along x-y directions ( $E_x$ ,  $E_y$ ,  $\nu_{xy}$ ,  $G_{xy}$ ) are defined as follows:

$$[\bar{S}] = \begin{bmatrix} \bar{S}_{11} & \bar{S}_{12} & \bar{S}_{16} \\ & \bar{S}_{22} & \bar{S}_{26} \\ Sym & & \bar{S}_{66} \end{bmatrix} = \begin{bmatrix} \frac{1}{E_x} & -\frac{\nu_{xy}}{E_x} & -\frac{m_x}{E_1} \\ & \frac{1}{E_y} & -\frac{m_y}{E_1} \\ Sym & & \frac{1}{G_{xy}} \end{bmatrix}$$

- $m_x$  and  $m_y$  are shear coupling parameters

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# ENGINEERING CONSTANTS FOR ANGLE LAMINA

$$[\bar{S}] = [T^*]^{-1}[S][T] \Rightarrow$$

$$\frac{1}{E_x} = \bar{S}_{11} = \frac{1}{E_1} c^4 + \left( \frac{1}{G_{12}} - \frac{2\nu_{12}}{E_1} \right) s^2 c^2 + \frac{1}{E_2} s^4$$

$$\frac{1}{E_y} = \bar{S}_{22} = \frac{1}{E_1} s^4 + \left( \frac{1}{G_{12}} - \frac{2\nu_{12}}{E_1} \right) s^2 c^2 + \frac{1}{E_2} c^4$$

$$\frac{1}{G_{xy}} = \bar{S}_{66} = 2 \left( \frac{2}{E_1} + \frac{2}{E_2} + \frac{4\nu_{12}}{E_1} - \frac{1}{G_{12}} \right) s^2 c^2 + \frac{1}{G_{12}} (s^4 + c^4)$$

$$\nu_{xy} = -E_x \bar{S}_{12} = E_x \left[ \frac{\nu_{12}}{E_1} (s^4 + c^4) - \left( \frac{1}{E_1} + \frac{1}{E_2} - \frac{1}{G_{12}} \right) s^2 c^2 \right]$$

$$m_x = -\bar{S}_{16} E_1 = -E_1 \left[ \left( \frac{-2}{E_1} - \frac{-2\nu_{12}}{E_1} - \frac{1}{G_{12}} \right) s c^3 - \left( \frac{2}{E_2} + \frac{2\nu_{12}}{E_1} - \frac{1}{G_{12}} \right) s^3 c \right]$$

$$m_y = -\bar{S}_{26} E_1 = -E_1 \left[ \left( \frac{-2}{E_1} - \frac{-2\nu_{12}}{E_1} - \frac{1}{G_{12}} \right) s^3 c - \left( \frac{2}{E_2} + \frac{2\nu_{12}}{E_1} - \frac{1}{G_{12}} \right) s c^3 \right]$$

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## HOOKE'S LAW FOR ANGLE LAMINA

- Up to now, we've expressed everything in terms of compliance, but we can express them in terms of stiffness too:

$$\begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix} = [Q] \begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix}$$

$$[T] \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = [Q][T^*] \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix}$$

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = [T]^{-1}[Q][T^*] \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix}$$

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = [\bar{Q}] \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = \begin{bmatrix} \bar{Q}_{11} & \bar{Q}_{12} & \bar{Q}_{16} \\ & \bar{Q}_{22} & \bar{Q}_{26} \\ Sym & & \bar{Q}_{66} \end{bmatrix} \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix}$$

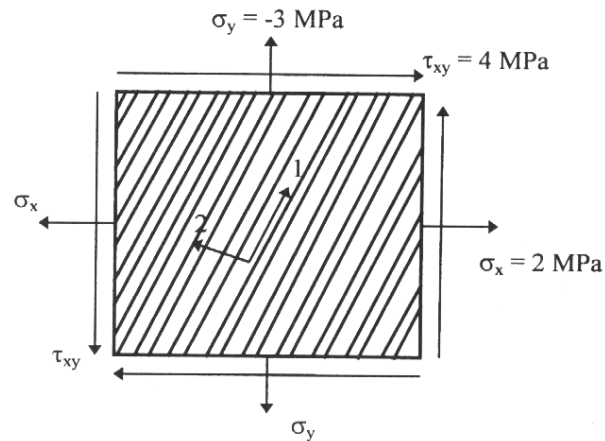
Note:  $[\bar{Q}]$  is referred to as the transformed stiffness matrix

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## ANGLE LAMINA EXAMPLE

- A 60° angle lamina is made from carbon/epoxy with properties as shown below
- With the applied stresses, determine:
  - the global strains
  - the local strains
  - the local stresses
  - the equivalent global engineering constants



$V_f$	0.70
$E_1$	181 GPa
$E_2$	10.30 GPa
$\nu_{12}$	0.28
$G_{12}$	7.17 GPa

## ANGLE LAMINA EXAMPLE SOLUTION

- First, find  $\cos q$ ,  $\sin q$

$$c = \cos(60^\circ) = 0.500$$

$$s = \sin(60^\circ) = 0.866$$

- Now, find the local compliance matrix

$$\begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix} = \begin{bmatrix} \frac{1}{E_1} & \frac{-\nu_{21}}{E_2} & 0 \\ \frac{-\nu_{12}}{E_1} & \frac{1}{E_2} & 0 \\ 0 & 0 & \frac{1}{G_{12}} \end{bmatrix} \begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix}$$

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## ANGLE LAMINA EXAMPLE SOLUTION

$$\left. \begin{aligned} S_{11} &= \frac{1}{181} = 5.52 \cdot 10^{-3} \text{ GPa}^{-1} \\ S_{12} &= -\frac{0.28}{181} = -1.55 \cdot 10^{-3} \text{ GPa}^{-1} \\ S_{22} &= \frac{1}{10.3} = 97.1 \cdot 10^{-3} \text{ GPa}^{-1} \\ S_{66} &= \frac{1}{7.17} = 139 \cdot 10^{-3} \text{ GPa}^{-1} \end{aligned} \right\} [S] = \begin{bmatrix} 5.52 & -1.55 & 0 \\ & 97.1 & 0 \\ \text{Sym} & & 139 \end{bmatrix} \cdot 10^{-3} \text{ GPa}^{-1}$$

And, the local stiffness matrix...

$$[Q] = [S]^{-1} = \begin{bmatrix} 181.8 & 2.897 & 0 \\ & 10.35 & 0 \\ \text{Sym} & & 7.17 \end{bmatrix} \text{ GPa}$$

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## ANGLE LAMINA EXAMPLE SOLUTION

$$[S] = \begin{bmatrix} 5.52 & -1.55 & 0 \\ & 97.1 & 0 \\ \text{Sym} & & 139 \end{bmatrix} \cdot 10^{-3} \text{ GPa}^{-1} \quad [\bar{S}] = [T^*]^{-1} [S] [T]$$

The global compliance matrix requires the transformation matrices...

$$[T] = \begin{bmatrix} c^2 & s^2 & 2sc \\ s^2 & c^2 & -2sc \\ -sc & sc & c^2 - s^2 \end{bmatrix} = \begin{bmatrix} 0.25 & 0.75 & 0.866 \\ 0.75 & 0.25 & -0.866 \\ -0.433 & 0.433 & -0.50 \end{bmatrix}$$

$$[T^*] = \begin{bmatrix} c^2 & s^2 & sc \\ s^2 & c^2 & -sc \\ -2sc & 2sc & c^2 - s^2 \end{bmatrix} = \begin{bmatrix} 0.25 & 0.75 & 0.433 \\ 0.75 & 0.25 & -0.433 \\ -0.866 & 0.866 & -0.50 \end{bmatrix}$$

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## ANGLE LAMINA EXAMPLE SOLUTION

And, the global compliance matrix becomes...

$$[\bar{S}] = [T^*]^{-1} [S] [T] = \begin{bmatrix} 80.53 & -7.878 & -32.34 \\ & 34.75 & -46.96 \\ \text{Sym} & & 114.1 \end{bmatrix} \cdot 10^{-3} \text{ GPa}^{-1}$$

The global stiffness matrix is the inverse...

$$[\bar{Q}] = [\bar{S}]^{-1} = \begin{bmatrix} 23.65 & 32.46 & 20.05 \\ & 109.4 & 54.19 \\ \text{Sym} & & 36.74 \end{bmatrix} \text{ GPa}$$

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## ANGLE LAMINA EXAMPLE SOLUTION

The engineering constants are found from the global compliance matrix...

$$\frac{1}{\bar{S}_{11}} = E_x = 12.4 \text{ GPa}$$

$$\nu_{xy} = -E_x \bar{S}_{12} = 0.0977$$

$$\frac{1}{\bar{S}_{22}} = E_y = 28.8 \text{ GPa}$$

$$m_x = -\bar{S}_{16} E_1 = 5.85$$

$$\frac{1}{\bar{S}_{66}} = G_{xy} = 8.76 \text{ GPa}$$

$$m_y = -\bar{S}_{26} E_1 = 8.50$$

The global strains are easily computed:

$$\begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = [\bar{S}] \begin{Bmatrix} 2 \\ -3 \\ 4 \end{Bmatrix} \text{ MPa} = \begin{Bmatrix} 55.34 \\ -307.8 \\ 532.8 \end{Bmatrix} \mu\varepsilon$$

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## ANGLE LAMINA EXAMPLE SOLUTION

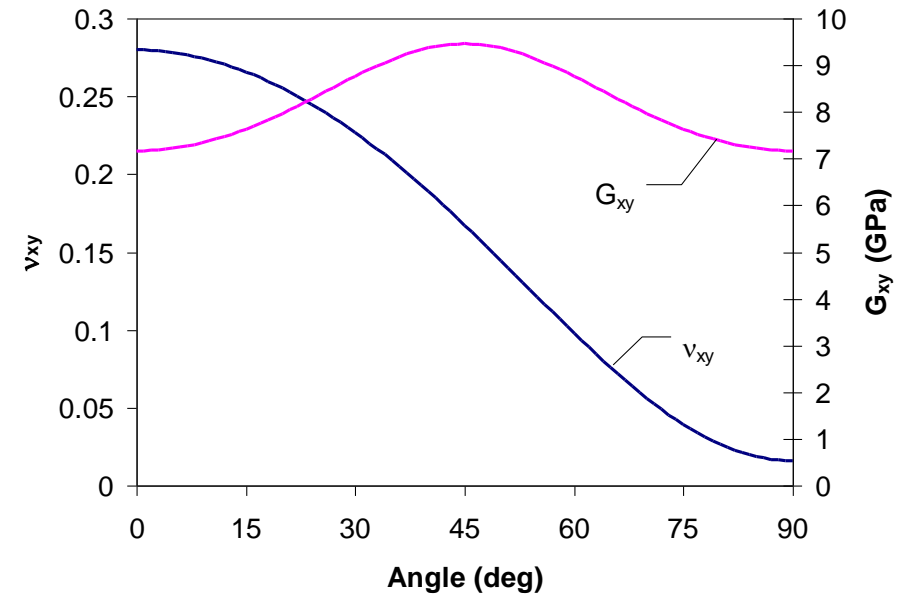
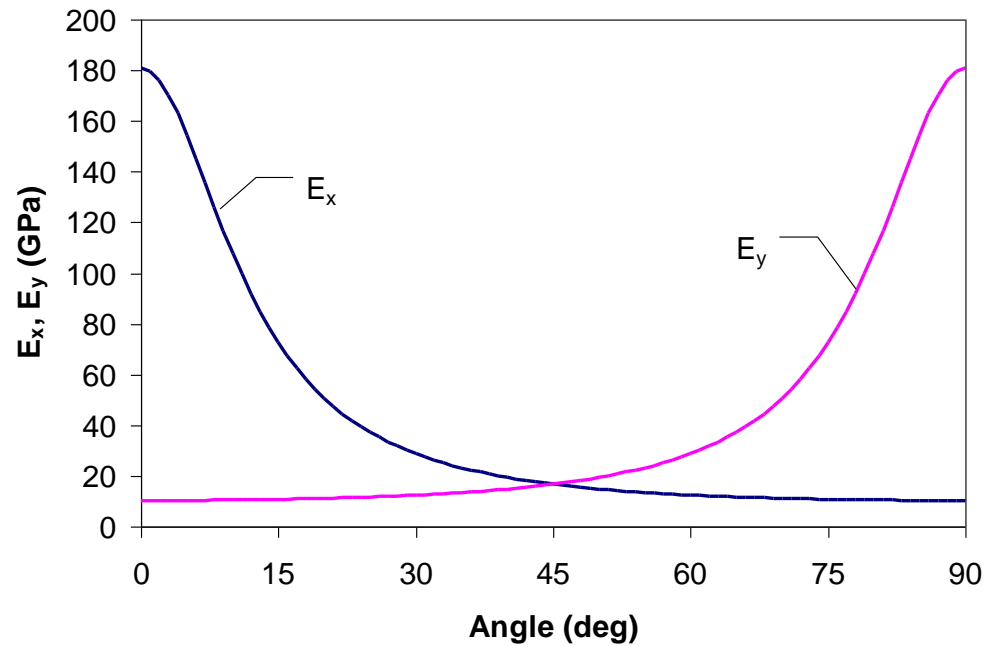
The local strains can be found by transforming the global strains:

$$\begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix} = [T^*] \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = \begin{Bmatrix} 13.67 \\ -266.2 \\ -580.9 \end{Bmatrix} \mu\varepsilon$$

The local stresses could be found by transforming the global stresses or by multiplying the local strains by the local stiffness matrix:

$$\begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix} = [Q] \begin{Bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{Bmatrix} = [T]^{-1} \begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = \begin{Bmatrix} 1.714 \\ -2.714 \\ -4.165 \end{Bmatrix} \text{MPa}$$

# ENGINEERING CONSTANTS AS A FUNCTION OF ANGLE



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**Thank you for joining us!**

The next session is:

***Session 8: Mechanics of Composites Part 2: Laminate Level***

*September 16, 2020 @ 9:00 am PT*

**Questions?**

For more information on future dates and times visit:

**[compositeskn.org](https://compositeskn.org)**

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